# Magnetocaloric effect in ordered double-perovskite Ba<sub>2</sub>FeMoO<sub>6</sub> synthesized using wet chemistry

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**Abstract.** Double perovskite  $Ba_2FeMoO_6$  has been prepared by a wet chemistry method and its magnetocaloric effect has been studied. The  $Ba_2FeMoO_6$  sample, consisting of crystallites with diameter of 130 nm and average aspect ratio of 15, exhibits a sharp magnetic transition. The maximum of the magnetic entropy change peaks at its Curie temperature (340 K), and a large magnetic entropy change (~1.54 J/kg K, or ~0.081 J/mol (atoms) K) is obtained upon 10 kOe applied magnetic field change. For comparison, the magnetic properties of  $Ba_2FeMoO_6$  prepared by a conventional solid-state reaction was also investigated.

**PACS.** 75.30.Sg Magnetocaloric effect, magnetic cooling – 82.80.-d Chemical analysis and related physical methods of analysis – 81.05.Mh Cermets, ceramic and refractory composites

## **1** Introduction

Half-metallic compounds, i.e. materials with only one spin direction present at the Fermi level, are interesting for their potential applications at room temperature in spin electronics. In this sense, oxides with high spin polarization and high Curie temperatures are desirable. Recently, double-perovskite oxides of the type  $A_2B'B''O_6$ , being also half-metallic ferromagnets, have received a great deal of attention because they display a substantial magnetoresistance and their Curie temperature is well above that of optimally doped simple manganite perovskites (such as  $La_{2/3}Ca_{1/3}MnO_3$  [1]. Usually, A is a large cation capable of 12-fold coordination with oxygen while B' and B''are different smaller cations suitable for octahedral coordination [2]. Among double perovskites, Fe-based oxides  $A_2$ FeMoO<sub>6</sub> (A =Sr, Ba, Ca, ...) have attracted much attention as a novel class of magnetoresistive oxides, since the low-field tunneling-type magnetoresistance effect was found in  $Sr_2FeMoO_6$  (SFMO) at room temperature [1]. It has been suggested that this compound has a double perovskite structure with alternately ordered  $\mathrm{Fe}^{3+}$  (3d<sup>5</sup>, S = 5/2) and Mo<sup>5+</sup> (4d<sup>1</sup>, S = 1/2) ions in the cubic  $ABO_3$  perovskite [3]. The Fe and Mo sublattices are ferromagnetically coupled within each sublattice, while the two sublattices are supposed to be antiferromagnetically

coupled. This gives rise to an S = 2 state and thus a saturation magnetization  $M_S = 4\mu_B \ (T = 0 \text{ K})$  is predicted. However, the  $M_S$  values of bulk materials reported so far are systematically much smaller than the predicted  $4\mu_B$ value. This difference between ideal and actual materials is believed to have its main origin in the lack of cationic ordering in the Fe/Mo sublattices. Since the ionic sizes of  $Fe^{3+}$  and  $Mo^{5+}$  are similar, there is a finite concentration of antisite defects in  $A_2$ FeMoO<sub>6</sub>, which interchange the positions of Fe and Mo in a random fashion. Each cation Fe that is misplaced in a Mo position is called an antisite (AS) defect. From the effects of AS defects on the saturation magnetization, it can be infered that since the sublattices Fe and Mo are antiferromagnetically arranged, each antisite contributes to a certain decrease of the saturation moment of the sample [4,5].

Many researchers have reported structural, magnetization, resistivity, Hall effect, specific heat and magnetic aftereffect of Fe-based double perovskite transition-metal oxides [6–20]. The properties of the double perovskites  $A_2$ FeMoO<sub>6</sub> (A =Ca, Sr and Ba) are thought to critically depend on the concentration of AS defects, which depends on the preparation method and synthesis conditions, such as heat-treatment temperature, time, and atmosphere. Most double perovskite materials reported previously were prepared by conventional solid-state reactions. Samples synthesized using conventional techniques frequently have

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isotropic grains with micrometer-scale size. Furthermore, to our knowledge, no studies have been carried on the magnetocaloric effect (MCE) of double perovskite transitionmetal oxides. In the present work, we synthesized double perovskite  $Ba_2FeMoO_6$  having high aspect ratio crystallites using wet chemistry, and studied its magnetocaloric effect above room temperature. For comparison, the structure, morphology and magnetic properties (including magnetocaloric effect) of  $Ba_2FeMoO_6$  prepared by a conventional solid-state reaction were also investigated.

#### 2 Experimental

Double perovskite Ba<sub>2</sub>FeMoO<sub>6</sub> was prepared by a wet chemistry method as follows: stoichiometric amounts of  $(NH_4)_6Mo_7O_{24} \cdot 4H_2O$ , Fe $(NO_3)_3 \cdot 9H_2O$  and BaCO<sub>3</sub> were dissolved in dilute HNO<sub>3</sub> solution at 333 K. Then suitable amounts of citric acid and ethylene glycol as coordinate agent were added in the reactor and a completely homogeneous light yellow transparent solution was achieved after heating at 343 K for 20 minutes. This solution was subjected to evaporation until a highly viscous residual was formed and a gel was developed during heating at 443 K. The gel was thermally treated at 773 K for 5 hours to decompose the organic precursor. After grinding, the sample was calcined in a stream of 7% H<sub>2</sub>/Ar at 1273 K for 5 h and furnace cooled.

In order to study the influence of preparation methods on the structure and properties of double perovskites,  $Ba_2FeMoO_6$  was also prepared by a conventional solidstate reaction method as follows: powders of  $BaCO_3$ ,  $Fe_2O_3$ , and  $MoO_3$  were mixed at stoichiometric ratios. Oxide mixtures were milled by high-energy ball milling for 5 h and then preheated at 873 K for 5 h under air. After further milling for another 5 h, the sample was sintered at 1273 K under an  $Ar/H_2$  (7%) atmosphere for 10 h. The process was repeated three times in order to complete the reaction.

It is noteworthy that in order to guarantee a precise sample composition, the sintering conditions must be carefully observed for both preparation methods. Higher temperatures and/or longer sintering times could produce different sample compositions because of the possibility of molybdenum oxide evaporation.

Phase identification and structural analysis were examined by X-ray powder diffraction (XRD) with Cu  $K\alpha$ radiation (Model D/Max-RA, Rigaku, Japan). The metal content in the sintered sample was determined by induced coupled plasma (ICP) spectroscopy (Model 1100 + 2000, Jarrell-Ash, USA). The oxygen content was also estimated by deducting the metal content from the weighed sample. The morphology of Ba<sub>2</sub>FeMoO<sub>6</sub> was examined directly by transmission electron microscopy [(TEM), Model JEM-200 CX, JEOL, Japan] and scanning electron microscopy [(SEM), JSM-5610 LV, JEOL, Japan]. Magnetization (M) versus temperature (T) and magnetization versus magnetic field (H) curves were measured by a vibrating sample magnetometer [(VSM), Model PAR 155, USA]. Curie temperature ( $T_C$ ) was determined from the



Fig. 1. XRD patterns of the two  $Ba_2FeMoO_6$  samples prepared by (a) wet chemistry and (b) solid-state reaction.

M-T curves. Magnetization of samples was measured in an isothermal regime under an applied magnetic field varying from 0 to 10 kOe. In the vicinity of  $T_C$ , isothermal M-H curves were obtained by a step of 3 or 5 K.

### 3 Results and discussion

The results of ICP indicate that the actual metal and oxygen content for both the samples prepared by wet chemistry, and the solid state reaction methods explored in this work, is practically equal to the nominal value. The X-ray diffraction patterns of the two samples prepared by wet chemistry and solid state reaction methods are shown in Figure 1. Both of them are single phase of  $Ba_2FeMoO_6$ without any detectable secondary phase, but the intensities of some peaks vary with preparation method. The diffraction peaks are indexed with respect to the cubic structure with a space group of Fm3m, where the unit cell has parameter a = 0.8071 (3) and a = 0.8078 (3) nm for samples prepared by wet chemistry and solid state reaction, respectively. These lattice parameters are similar to the value (0.80747 nm) reported by Kim et al. [9] but smaller than the value (0.8062 nm) obtained by Borges et al. [19] The concentration of antisite (AS) defects of the Fe and Mo arrangement were obtained from Rietveld refinements. The AS concentration is defined as the ratio of the concentration of Fe cations in Mo positions divided by the total Fe cations. Therefore, AS = 50% corresponds to a totally disordered double perovskite. The AS defect concentrations were estimated to be 2% and 9% for samples prepared by wet chemistry and solid state reaction, respectively. For the wet chemistry method used in our

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Fig. 2. Representative TEM images at two different magnifications of the  $Ba_2FeMoO_6$  prepared by wet chemistry.

experiment, the reactant is a heteropolymer of molybdates. When the reactant is preheated, it becomes a mixture of metal oxides, which are fully uniformly mixed at the atomic scale. Thus it is easier to transform into  $Ba_2FeMoO_6$  with a high order parameter.

Figures 2a and b show representative TEM images of  $Ba_2FeMoO_6$  prepared by wet chemistry. One can see that the sample exhibits bullrush-like crystallites with diameter of 130 nm and high aspect ratio larger than 15. Figure 3 presents a typical micrograph of  $Ba_2FeMoO_6$  prepared by solid state reaction. One can observe that the sample consists of inhomogeneous aggregates of granules with different shape and size.

The ferromagnetic ordering transition temperature, defined as the temperature at which the  $dM/dT \sim T$  curve reaches a minimum, has been extracted from the low field (1 kOe) magnetization versus temperature. The  $T_C$  values vary little for the two samples. Figure 4 displays the temperature dependence of magnetization M for the two samples Ba<sub>2</sub>FeMoO<sub>6</sub>. The magnetization was measured upon warming in a 10 kOe field after cooling to 100 K in zero field. The M at 173 K in a 10 kOe field is 37.3 emu/g  $(3.49 \ \mu_B/\text{per formula unit})$  for Ba<sub>2</sub>FeMoO<sub>6</sub> prepared by wet chemistry. As can be seen, the magnetization of the sample prepared by solid state reaction is lower than that of the sample prepared by wet chemistry. Further investigations into the magnetic behavior at lower temperature are in progress. Although the magnetization at 173 K is not saturated, the tendency of saturation magneticza-



Fig. 3. SEM image of the  $Ba_2FeMoO_6$  prepared by solid-state reaction.



Fig. 4. Temperature dependence of magnetization (warming curves) for zero-field cooled (ZFC)  $Ba_2FeMoO_6$ . Hollow circles: sample prepared by wet chemistry; solid squares: sample prepared by solid-state reaction.

tion  $(M_S)$  may be estimated from Figure 4. This decrease of magnetization would be due to antisite defects resulting from the partial disorder of Fe and Mo ions among the B'/B'' sublattices. Balcells et al. [2] observed that the experimental saturation magnetization had an almost linear dependence on the AS concentration. Monte Carlo calculations of Ogale et al. [4] have also predicted a reduction of magnetization as a function of the AS defect concentration:

$$M_S(\mu_B/f.u.) = 4 - 8 \times (AS).$$
 (1)

Sakuma et al. [21] discussed the dependence of the spontaneous magnetization  $(M_0)$  on the order parameter (x) using two models considering the local magnetic interaction between ions on the B sites, and reported the relationships between  $M_0$  and x, respectively, as

$$M_0(\mu_B/\text{f.u.}) = x/25$$
 (2)

and

$$M_0(\mu_B/\text{f.u.}) = x/20 - 1.$$
 (3)

The nature and origin of this decrease of the magnetisation in the presence of disorder is still a matter of debate in the literature. There are two distinct ways that the net magnetisation may be reduced in the presence of antisite defects [14]. One possibility is that the disorder destroys the specific spin arrangement of Fe and Mo sublattices without any significant effect on the individual magnetic moments at these sites. This can be achieved by transforming the ferromagnetic coupling between some of the Fe sites to an antiferromagnetic coupling. Alternately, the magnetic moments at each individual site may decrease due to the different chemical environments induced by the disorder, without affecting the nature of the spin order within the Fe and Mo sublattices. The real situation may even be a combination of both these effects, with a simultaneous reduction in the magnetic moments at different sites as well as a change in the nature of the magnetic coupling between different sites.

Based on the thermodynamic theory, the entropy change, which results from the spin ordering (i.e., ferromagnetic ordering) and is induced by the variation of the applied magnetic field from 0 to  $H_{max}$ , is given by [22,23]

$$|\Delta S_M| = \int_0^{H_{\text{max}}} \left(\frac{\partial M}{\partial T}\right) dH.$$
 (4)

According to equation (4), the magnetic entropy change  $(|\Delta S_M|)$  depends on the temperature gradient of the magnetization and attains a maximum value around the Curie temperature, at which the magnetization decays most rapidly. The big advantage of the magnetization measurements is the fact that they can be carried out rather quickly, and thus they can be used to evaluate the magnetocaloric properties of potential magnetic refrigerant materials. Isothermal magnetization curves of Ba<sub>2</sub>FeMoO<sub>6</sub> have been measured at various temperatures with a vibrating sample magnetometer (VSM). Figure 5 shows a part of the isothermal magnetization curves for the Ba<sub>2</sub>FeMoO<sub>6</sub> prepared by wet chemistry. In our experiments, the rate of changing the magnetic field (20 Oe/s)is slow enough to get an isothermal  $M \sim T$  curve. In Figure 6, we show the temperature dependence of magnetic entropy change  $(|\Delta S_M|)$  under an applied field of 10 kOe for Ba<sub>2</sub>FeMoO<sub>6</sub>. Upon 10 kOe applied field change, the maximum  $|\Delta S_M|$  of Ba<sub>2</sub>FeMoO<sub>6</sub> prepared by wet chemistry and solid state reaction is 1.54 and 1.09 J/kgK, respectively. The peak of the  $|\Delta S_M|$  for the sample prepared by solid state reaction becomes broader. Since the molecular weight of Ba<sub>2</sub>FeMoO<sub>6</sub> is very large compared with that of  $ABO_3$ -type perovskites, the magnetic entropy change per mole of atoms  $(|\Delta S_M|_{mol})$  of Ba<sub>2</sub>FeMoO<sub>6</sub> is rather large (0.081 and 0.057 J/mol (atoms) K). For comparison, we summarize in Table 1 the data of several typical perovskite materials that order magnetically near room temperature (>250 K).



Fig. 5. Magnetization as a function of applied field measured at different temperatures for the  $Ba_2FeMoO_6$  prepared by wet chemistry. Only a part of isothermal magnetization cures are presented.



Fig. 6. Temperature dependence of magnetic entropy change in an applied field  $H_{max} = 10$  kOe for Ba<sub>2</sub>FeMoO<sub>6</sub>. Hollow circles: sample prepared by wet chemistry; solid squares: sample prepared by solid-state reaction.

According to relation (4), a large magnetocaloric effect can be expected for materials that exhibit a sharp magnetic transition. Herein, we prepared the Ba<sub>2</sub>FeMoO<sub>6</sub> using wet chemistry and synthesized at thermodynamic conditions assuring nearly ideal stoichiometry of the material with a high order parameter of the Fe and Mo arrangement. However, sample synthesized by solid state reaction showed inhomogeneous aggregates of granules with a high AS defect concentration, which negatively affects the sharpness of the magnetic transition and gives rise to a low  $|\Delta S_M|$ .

#### 4 Conclusions

The structure, morphology and magnetic properties (including magnetocaloric effect) of double perovskite  $Ba_2FeMoO_6$  prepared by wet chemistry have been compared with those of sample synthesized by solid state

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Sample	$T_C$	$\Delta H$	$ \Delta S_M $	$^{\mathrm{a}} \Delta S_{M} _{mol}$	Source of
•	(K)	$(\times 10 \text{ kOe})$	(J/kg K)	(J/mol (atoms) K)	Data
Ba <sub>2</sub> FeMoO <sub>6</sub> <sup>b</sup>	340	1.0	1.54	0.081	Present work
$\operatorname{Ba_2FeMoO_6}^{\operatorname{c}}$	340	1.0	1.09	0.057	Present work
$La_{2.061}K_{0.935}Mn_{1.997}O_7$	254	1.0	1.1	0.050	[24]
$La_{2/3}Ba_{1/3}MnO_{3.00}$	337	1.0	2.77	0.134	[25]
${\rm La}_{0.834} {\rm Na}_{0.163} {\rm Mn}_{1.000} {\rm O}_3$	343	1.0	2.11	0.094	[26]
$\rm La_{0.796}K_{0.196}Mn_{0.993}O_{3}$	344	1.0	1.55	0.068	[26]
$La_{0.75}Sr_{0.15}Ca_{0.10}MnO_{3}$	327	1.5	2.8	0.126	[27]
$La_{0.75}Sr_{0.125}Ca_{0.125}MnO_{3}$	283	1.5	1.5	0.068	[27]
$La_{0.65}Ca_{0.35}MnO_3$	$\sim 255$	1.0	$\sim 2.7$	0.112	[28]
$La_{0.67}Ca_{0.33}MnO_{\delta}$	260	1.0	1.1	0.046	[29]

Table 1. Room-temperature magnetic entropy change  $(|\Delta S_M| \text{ and } |\Delta S_M|_{mol})$  of perovskite-type materials at their Curie temperature (>250 K) in low magnetic fields.

<sup>a</sup> molar magnetic entropy change calculated from  $|\Delta S_M|$ ;

<sup>b</sup> prepared by wet chemistry;

<sup>c</sup> prepared by solid state reaction.

reaction. The use of wet chemistry method enables the formation of double perovskite  $Ba_2FeMoO_6$  containing high aspect ratio crystallites with a high order parameter of the Fe and Mo arrangement. The sample prepared by wet chemistry exhibits a large magnetization, a sharp magnetic transition and a large magnetic entropy change under a low magnetic field above room temperature. Although the properties of our materials are far from optimal, further exploration of this double perovskite may prove fruitful.

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#### References

- K.I. Kobayashi, T. Kimura, H. Sawada, K. Terakura, Y. Tokura, Nature **395**, 677 (1998)
- LI. Balcells, J. Navarro, M. Bibes, A. Roig, B. Martínez, J. Fontcuberta, Appl. Phys. Lett. 78, 781 (2001)
- M. Itoh, I. Ohta, Y. Inaguma, Mater. Sci. Eng. B 41, 55 (1996)
- A.S. Ogale, S.B. Ogale, R. Ramesh, T. Venkatesan, Appl. Phys. Lett. **75**, 537 (1999)
- J. Navarro, J. Nogue's, J.S. Muñoz, J. Fontcuberta, Phys. Rev. B 67, 174416 (2003)
- Y. Moritomo, S. Xu, A. Machida, T. Akimoto, E. Nishibori, M. Takata, M. Sakata, Phys. Rev. B 61, R7827 (2000)
- J. Navarro, C. Frontera, L.I. Balcells, B. Martínez, J. Fontcuberta, Phys. Rev. B 64, 092411 (2001)
- J. Navarro, LI. Balcells, F. Sandiumenge, M. Bibes, A. Roig, B. Martínez, J. Fontcuberta, J. Phys.: Condens. Mater 13, 8481 (2001)
- S.B. Kim, B.W. Lee, S.R. Yoon, C.S. Kim, J. Magn. Magn. Mater. 254–255, 580 (2003)

- D. Niebieskikwiat, A. Caneiro, R.D. Sanchez, J. Fontcuberta, Phys. Rev. B 64, 180406 (2001)
- B.G. Kim, Y.S. Hor, S.W. Cheong, Appl. Phys. Lett. 79, 388 (2001)
- D.D. Sarma, P. Mahadevan, T. Saha-Dasgupta, S. Ray, A. Kumar, Phys. Rev. Lett. 85, 2549 (2000)
- J. Linden, T. Yamamoto, M. Karppinen, H. Yamauchi, T. Pietari, Appl. Phys. Lett. 76, 2925 (2000)
- D.D. Sarma, Current Opinion in Solid State and Mater. Sci. 5, 261 (2001)
- H.Yanagihara, W. Cheong, M.B. Salamon, S. Xu, Y. Moritomo, Phys. Rev. B 65, 092411 (2002)
- G. Popov, M. Greenblatt, M. Croft, Phys. Rev. B 67, 024406 (2003)
- C.L. Yuan, Y. Zhu, P.P. Ong, Appl. Phys. Lett. 82, 934 (2003)
- K. Phillips, A. Chattopadhyay, A.J. Millis, Phys. Rev. B 67, 125119 (2003)
- R.P. Borges, R.M. Thomas, C. Cullinan, J.M.D. Coey, R. Suryanarayanan, L. Ben-Dort, L. Pinsard-Gaudart, A. Revcolevschi, J. Phys.: Condens. Mater 11, L445 (1999)
- Y. Tomioka, T. Okuda, Y. Okimoto, R. Kumai, K.I. Kobayashi, Y. Tokura, Phys. Rev. B 61, 422 (2000)
- H. Sakuma, T. taniyama, Y. Kitamoto, Y. Yamazaki, J. Appl. Phys. 93, 2816 (2003)
- V.K. Pecharsky, K.A. Gschneidner Jr, J. Magn. Magn. Mater. 200, 44 (1999)
- V.K. Pecharsky, K.A. Gschneidner Jr, J. Appl. Phys. 90, 4614 (2001)
- W. Zhong, W. Chen, H.Y. Jiang, X.S. Liu, C.T. Au, Y.W. Du, Eur. Phys. J. B **30**, 331 (2002)
- W. Zhong, W. Chen, C.T. Au, Y.W. Du, J. Magn. Magn. Mater. 261, 238 (2003)
- W. Zhong, W. Chen, W.P. Ding, N. Zhang, A. Hu, Y.W. Du, Q.J. Yan, Eur. Phys. J. B 3, 169 (1998)
- 27. Z.B. Guo, Y.W. Du, J.S. Zhu, H. Huang, W.P. Ding, D. Feng, Phys. Rev. Lett. 78, 1142 (1997)
- X. Bohigas, E. del Barco, M. Sales, J. Tejada, J. Magn. Magn. Mater. **196**, 455 (1999)
- X.X. Zhang, J. Tejada, Y. Xin, G.F. Sun, K.W. Wong, X. Bohigas, Appl. Phys. Lett. 69, 3596 (1996)